

## Low-lying structure of $^{50}\text{Ar}$ and the $N = 32$ subshell closure<sup>†</sup>

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It is now well known that far from the line of  $\beta$  stability the nuclear magic numbers can change from their standard values. For example, in the  $pf$  shell, the onset of a new magic number at  $N = 32$  has been reported along the Cr, Ti, and Ca isotopic chains, while a sizable gap at  $N = 34$  was deduced from the structure of  $^{54}\text{Ca}^{1)}$ . Very recently, the persistence of the  $N = 32$  subshell closure was established in systems below the  $Z = 20$  core<sup>2)</sup>. In the present work, the low-lying structure of  $^{50}\text{Ar}$  has been investigated to shed light on the character of the  $N = 32$  magic number at more extreme neutron-to-proton ratios. Preliminary results are discussed in Ref.<sup>3)</sup>.

A primary beam of  $^{70}\text{Zn}^{30+}$  ions with a typical intensity of  $\sim 60$  pnA was used to generate a fast radioactive beam containing  $^{54}\text{Ca}$ ,  $^{55}\text{Sc}$ , and  $^{56}\text{Ti}$ , amongst other products. The constituents were identified using the BigRIPS separator and focused on a 10-mm-thick  $^9\text{Be}$  reaction target at the eighth focal plane. Reaction products emerging from the target were identified by the ZeroDegree spectrometer (ZDS); despite the fact that the ZDS was optimized for the transmission of  $^{54}\text{Ca}^{1)}$ , a sufficient number of  $^{50}\text{Ar}$  ions fell within the acceptance of the spectrometer to extract structural information. The reaction target was surrounded by the DALI2  $\gamma$ -ray detector array to measure transitions emitted from nuclear excited states.

The  $\gamma$ -ray energy spectra—corrected for the Doppler effect—are presented in Fig. 1 using the sum of the  $^9\text{Be}(^{54}\text{Ca}, ^{50}\text{Ar})X$ ,  $^9\text{Be}(^{55}\text{Sc}, ^{50}\text{Ar})X$ , and  $^9\text{Be}(^{56}\text{Ti}, ^{50}\text{Ar})X$  multinucleon removal reactions. The line at 1178(18) keV, which is the most intense peak in the spectra, is assigned as the transition from the yrast  $2^+$  state to the  $0^+$  ground state in  $^{50}\text{Ar}$ . A weaker, tentative peak is present at 1582(38) keV, and is suggested as the transition between the  $4_1^+$  and  $2_1^+$  levels. Statistics were insufficient to confirm the proposed decay scheme using  $\gamma\gamma$  coincidence relationships.

The  $2_1^+$  state in  $^{50}\text{Ar}$  indicates an increase in energy relative to its even-even neighbour  $^{48}\text{Ar}$  and, therefore, naively suggests the presence of a sizable subshell closure at  $N = 32$  in Ar isotopes. In order to investigate the nature of the increase in energy in more detail, large-scale shell-model calculations employing a modified version of the SDPF-MU Hamiltonian<sup>4)</sup> were performed; the modifications were based on recent experimental data from exotic  $\text{Ca}^{1)}$  and  $\text{K}^{5)}$  isotopes. The predictions reproduce the experimental energy levels in lighter Ar isotopes, and the results of the present work, in a satisfactory manner. Moreover, the calculations indicate that the magnitude of the  $N = 32$  subshell closure in  $^{50}\text{Ar}$  is equally as significant as the gaps in  $^{52}\text{Ca}$  and  $^{54}\text{Ti}$ , where the experimental evidence for this magic number is well documented. The calculations also indicate a rather high  $2_1^+$  energy in  $^{52}\text{Ar}$  and, therefore, experimental input on this nucleus is encouraged to investigate the significance of the  $N = 34$  subshell closure in more exotic systems.

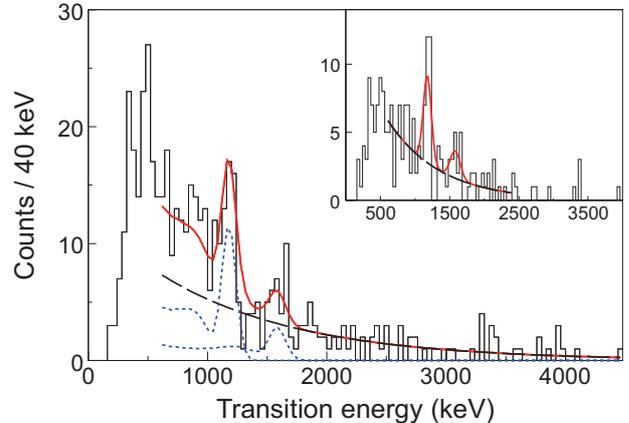


Fig. 1. (colour) Doppler-corrected  $\gamma$ -ray energy spectra for  $^{50}\text{Ar}$ . The main and inset panels display  $M_\gamma \geq 1$  and  $M_\gamma \leq 3$  data, respectively. The black dashed lines are exponential fits to the backgrounds and the blue dashed lines are GEANT4 simulations; the solid red lines are total (sum) fits.

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