

# Successive Transitions in Spin-dimer Compound $\text{Cs}_3\text{V}_2\text{Cl}_9$

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$\text{A}_3\text{M}_2\text{X}_9$  (A = Cs, Rb : M = transition metal elements: X = Cl, Br) compounds with trigonal space group  $\text{P6}_3/\text{mmc}$  are composed of isolated di-nuclear complexes  $[\text{M}_2\text{X}_9]^{3-}$  and their magnetic properties are explained within an isolated or weakly coupled spin dimer model. Because these spin dimers are arranged in a triangular form, the spin frustration effect is expected to appear when magnetic phase transition occurs via finite interdimer interactions.  $\text{Cs}_3\text{V}_2\text{Cl}_9$ , one member of the  $\text{A}_3\text{M}_2\text{X}_9$  family with magnetic ion  $\text{V}^{3+}$  ( $S = 1$ ), was previously studied via magnetic susceptibility and inelastic neutron scattering measurement using powder sample<sup>1)</sup> and no magnetic ordering was observed above 1.5 K. We recently synthesized a single crystal of  $\text{Cs}_3\text{V}_2\text{Cl}_9$  and measured susceptibility  $\chi(T)$  and specific heat, and we found successive phase transition at  $T_N \approx 4$  K and  $T_n \approx 15$  K.<sup>2)</sup>

These successive phase transitions are not explained within the framework of the isolated dimer model, and they show the presence of non-negligible interdimer interaction.  $\chi(T)$  shows no anomaly at  $T_n$  although  $T_N$  is accompanied by the cusp-like anomaly of  $\chi(T)$ . The lower temperature transition is suggested to be an antiferromagnetic transition although the spin structure is not clear. The higher temperature transition is not a mere crystal structure transition because the transition temperature depends on the applied magnetic field.

We measured the  $\mu\text{SR}$  of  $\text{Cs}_3\text{V}_2\text{Cl}_9$  to clarify the nature of the successive phase transitions. Figure 1 shows the zero-field muon spin relaxation (ZF- $\mu\text{SR}$ ) spectra

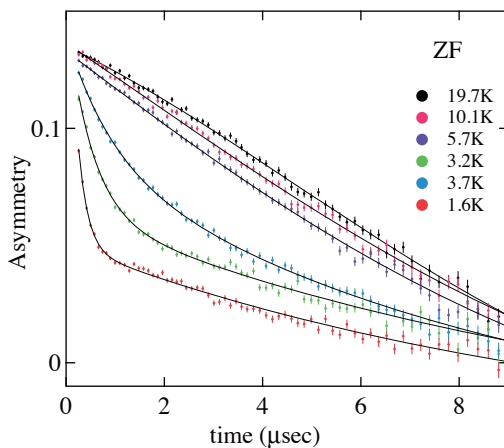


Fig. 1. Temperature dependence of the ZF- $\mu\text{SR}$  spectra of  $\text{Cs}_3\text{V}_2\text{Cl}_9$ .

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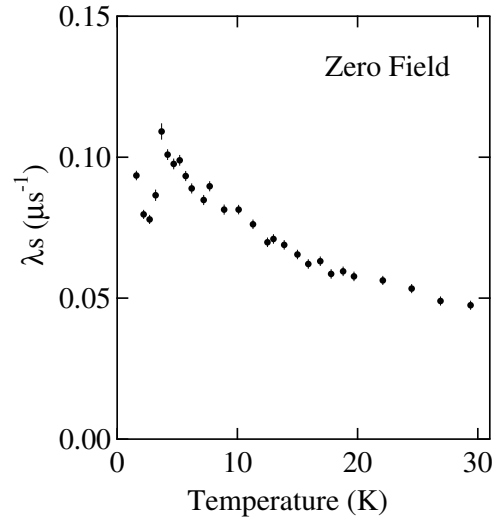


Fig. 2. Temperature dependence of the relaxation rate  $\lambda_s$ .

measured down to 1.6 K. At relatively high temperatures, the spectra follow Gaussian curves. As the temperature decreases, the spectra changes from Gaussian to an exponential curve following  $a_s \exp(-\lambda_s t) + a_f \exp(-\lambda_f t)$ , where  $\lambda_s$  and  $\lambda_f$  denote slow and fast relaxation rates, and  $a_s$  and  $a_f$  represent amplitudes of the asymmetry of slow and fast components, respectively. Solid lines in Fig. 1 are fitted results. Figure 2 shows the temperature dependence of  $\lambda_s$ . A distinct peak is observed at around  $T_N$ , which confirms that a magnetic long range order occurs at this temperature. No anomaly of  $\lambda_s$  is observed at  $T_n$ , which indicates that an internal field does not appear. The phase transition at  $T_n$  is not accompanied by the internal field, whereas the value of  $T_n$  depends on the applied magnetic field. One candidate for the transition at  $T_n$  is a spin nematic order wherein quadrupole moments, not magnetic moments, play the role of an order parameter. Because the nematic state does not break the time-reversal symmetry, usual magnetic probes including muon do not detect this transition.<sup>3)</sup> Several theoretical studies indicate the occurrence of the nematic order in an  $S = 1$  triangular lattice antiferromagnet or spin dimer magnets. The findings obtained by this  $\mu\text{SR}$  experiments make it more likely that  $T_n$  is the nematic transition.

## References

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