

Dineutron and effective pairing forces in momentum space

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The spatial two-neutron correlation, called dineutron correlation, is one of the unique features around the neutron drip line. The dineutron correlation has been discussed extensively in light-mass nuclei such as ^{11}Li ,^{1,2)} but it is considered to be a universal phenomenon over all mass-number regions.³⁾

We propose a new effective pairing force by focusing on the spatial structure of a neutron pair. First, we discuss the necessity by performing a three-body model calculation with the density-dependent zero-range force (DD $\delta(\mathbf{r})$ -force),^{1,2)} $V_\delta(\mathbf{r}_1, \mathbf{r}_2) = V^{(0)} \{1 - \eta[\rho(\mathbf{r}_1)/\rho_0]\} \delta(\mathbf{r}_1 - \mathbf{r}_2)$. The DD $\delta(\mathbf{r})$ -force is widely used in nuclear structure calculations.¹⁻³⁾ It must be supplemented with a cutoff in the two-particle spectrum, $\varepsilon_1 + \varepsilon_2 \leq E_{\text{cut}}$. The parameters $V^{(0)}$ and η are adjusted for each E_{cut} so as to reproduce the properties of low-energy neutron-neutron (nn) scattering and the two-neutron energy E_{2n} of finite nuclei. The parameter $\rho_0 = 0.16 \text{ fm}^{-3}$ is the saturation density.

In Fig. 1, the root-mean-square (rms) values k_{rel} and q_{cm} of the relative and center of mass (cm) momenta of paired neutrons in ^{11}Li and ^{12}Be are shown. k_{rel} converges, while q_{cm} diverges as a function of E_{cut} owing to undesirable coupling to high-momentum components of single-particle states in the continuum.

In order to overcome the difficulties, we consider a new effective pairing force in momentum space (k -SEP force), $V_{\text{SEP}} = -\frac{1}{2\pi^2} V_k^{(0)} g(\mathbf{k})g(\mathbf{k}')h(\mathbf{q})h(\mathbf{q}')$. Here, the Yamaguchi-type form factor $g(\mathbf{k}) = 1/(\mathbf{k}^2 + \Lambda^2)$ and the Gaussian distribution $h(\mathbf{q}) = (\sqrt{\pi}q_0)^{-3} e^{-\mathbf{q}^2/q_0^2}$ for the relative momentum \mathbf{k} and the cm momentum \mathbf{q} of paired neutrons are adopted. The parameters $V_k^{(0)}$ and

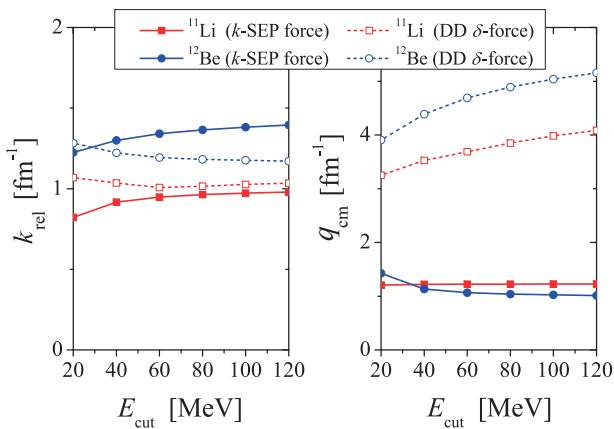


Fig. 1. Relative momentum k_{rel} and cm momentum q_{cm} of the neutron pair in ^{11}Li and ^{12}Be . The results using the k -SEP force and the DD $\delta(\mathbf{r})$ -force are shown.

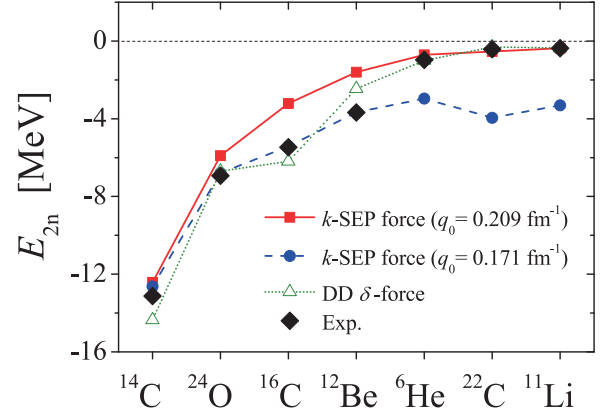


Fig. 2. Two-neutron energies E_{2n} obtained with the k -SEP force ($q_0 = 0.209$ and 0.171 fm^{-1}) and the DD $\delta(\mathbf{r})$ -force with fixed $E_{\text{cut}} = 40 \text{ MeV}$ and $\eta = 0.849$, in comparison with the experimental data.

Λ are fixed so as to reproduce the properties of low-energy nn scattering.⁴⁾ The parameter q_0 is fixed by the two-neutron energy E_{2n} of finite nuclei.

The k_{rel} and q_{cm} obtained with the k -SEP force are shown in Fig. 1. They converge well above $E_{\text{cut}} = 100 \text{ MeV}$. Here, the parameter $q_0 = 0.209 \text{ fm}^{-1}$ (0.171 fm^{-1}) is used for ^{11}Li (^{12}Be). The small cm momentum q_{cm} indicates that the pairing correlation occurs only around the nuclear surface region in real space. The relative momentum k_{rel} in ^{12}Be becomes larger than that in Borromean ^{11}Li owing to the high-momentum component in the bound $1p_{1/2}$ state.

The two-neutron energies E_{2n} of three-body systems can be classified into two categories as shown in Fig. 2. The Borromean nuclei (^6He , ^{22}C , ^{11}Li) are well described with $q_0 = 0.209 \text{ fm}^{-1}$, while $q_0 = 0.171 \text{ fm}^{-1}$ for non Borromean nuclei. Because the parameter q_0 introduces a cutoff for correlations in the opening angle θ_k between \mathbf{k}_1 and \mathbf{k}_2 as $\cos \theta_k \lesssim q_0^2/(2k_1k_2)$ in $h(\mathbf{q})$, the pairing correlations have a more collective nature in calculation with a larger q_0 .

In conclusion, we proposed a new effective pairing force that describes well the structure of the neutron pair. It is also easily utilizable in various frameworks such as the density functional theory in k space.⁵⁾

References

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